

## CHAPTER 7: THE DE FINETTI THEOREM

### § 7.1 Extendibility of quantum states

Consider a bipartite state  $\rho_{AB}$ . We call  $\rho_{AB}$  *k-extendible* if there exists a *k-extension*, a state  $\rho_{AB_1 \dots B_k}$  with  $B_i \cong B$  and  $\rho_{AB_i} = \text{tr}_{B_1 \dots B_{i-1} B_{i+1} \dots B_k} \rho_{AB_1 \dots B_k} = \rho_{AB}$  for all  $i=1, \dots, k$ .

**Lem** Separable states are  $\infty$ -extendible.

Proof: Let  $\rho_{AB} = \sum_i p_i \sigma_A^{(i)} \otimes \sigma_B^{(i)}$  be separable, then

$$\rho_{AB_1 \dots B_k} = \sum_i p_i \sigma_A^{(i)} \otimes \sigma_{B_1}^{(i)} \otimes \dots \otimes \sigma_{B_k}^{(i)}$$
 defines a

*k-extension* for arbitrary  $k \in \mathbb{N}$ . □

Conversely, one can show that for every entangled state  $\rho_{AB}$  there exists a  $k_0$  such that  $\rho_{AB}$  has no *k-extension* for  $k \geq k_0$ .

Ex.: Pure entangled states are not even 2-extendible.

This is usually called *monogamy of entanglement*:

A quantum system cannot be entangled with a large number of other systems.

*De Finetti theorems* provide a quantitative version of monogamy.

## § 7.2 A De Finetti theorem for pure symmetric states

We will focus on pure states in the symmetric subspace:

$$|\psi\rangle \in \text{Sym}^n(\mathbb{C}^d) = \{ |\psi\rangle \in (\mathbb{C}^d)^{\otimes n} : \varphi(\pi) |\psi\rangle = |\psi\rangle \forall \pi \in S_n \}$$

Note that  $\dim \text{Sym}^n(\mathbb{C}^d) = \binom{n+d-1}{n}$  by Weyl's dim. formula

**Lem** Let  $|\varphi\rangle \in \mathbb{C}^d$  be arbitrary. Then

$$\pi_n = \binom{n+d-1}{n} \int_{\mathcal{U}_d} dU (U |\varphi\rangle \langle \varphi| U^\dagger)^{\otimes n}$$

is equal to the projector onto the symmetric subspace.

Proof: This follows from showing that

a)  $\pi_n \in \text{End}(\text{Sym}^n(\mathbb{C}^d))$

b)  $\mathcal{U}_d$  acts irreducibly on  $\text{Sym}^n(\mathbb{C}^d)$  via  $U \mapsto U^{\otimes n}$

c)  $\pi_n$  is invariant under this action and Schur's lemma.

Details are left as an exercise. □

Recall the trace distance  $D(\rho, \sigma) := \frac{1}{2} \|\rho - \sigma\|_1$

(with the trace norm  $\|x\|_1 = \text{tr} \sqrt{x+x^\dagger}$ ),

and the fidelity  $F(\rho, \sigma) = \|\sqrt{\rho} \sqrt{\sigma}\|_1$ , and  $D \leq \sqrt{1 - F^2}$ .

## Prop (De Finetti theorem for pure symmetric states)

Let  $\mathcal{H}_{A_i} \cong \mathbb{C}^d$  and  $|\psi\rangle_{A_1 \dots A_n} \in \text{Sym}^n(\mathbb{C}^d)$ .

Then for any  $k < n$ ,

$$D\left(\rho_{A_1 \dots A_k}, \int d\psi p(\psi) |\psi\rangle\langle\psi|^{\otimes k}\right) \leq \sqrt{\frac{dk}{n-k}}.$$

Here,  $d\psi$  is the measure on pure states induced by the Haar measure, i.e.,  $|\psi\rangle = U|\psi_0\rangle$  for some fixed  $|\psi_0\rangle$  and  $d\psi f(|\psi\rangle) = dU f(U|\psi_0\rangle)$ .

$p(\psi)$  is a probability density that depends on  $|\psi\rangle$ .

Proof: The main idea is the following: For  $m = n - k$  interpret  $\Pi_m = \binom{m+d+1}{m} \int d\psi |\psi\rangle\langle\psi|^{\otimes m}$  as a continuous PDM with operators  $\binom{m+d+1}{m} |\psi\rangle\langle\psi|^{\otimes m}$  measuring the last  $m$  systems that we trace out. Getting a specific outcome  $|\psi\rangle \in \mathbb{C}^d$  should then imply that the first  $k$  systems are also in the state  $|\psi\rangle^{\otimes k}$  on average, due to the permutation invariance of  $|\psi\rangle_{A_1 \dots A_n}$ .

Since  $|\psi\rangle_{A_1 \dots A_n}$  is symmetric under arbitrary perm.s,  
 in particular  $|\psi\rangle_{A_1 \dots A_n} = (\mathbb{1}_k \otimes \Pi_m) |\psi\rangle_{A_1 \dots A_n}$   
 (recall  $n = k+m$ ). Hence,

$$\begin{aligned} \chi_{A_1 \dots A_k} &= \text{tr}_{A_{k+1} \dots A_n} |\psi\rangle\langle\psi|_{A_1 \dots A_n} \\ &= \text{tr}_{A_{k+1} \dots A_n} \left[ (\mathbb{1}_k \otimes \Pi_m) |\psi\rangle\langle\psi|_{A_1 \dots A_n} \right] \\ &= \binom{m+d-1}{m} \int d\varphi (\mathbb{1}_k \otimes \langle\varphi|^{\otimes m}) |\psi\rangle\langle\psi| (\mathbb{1}_k \otimes |\varphi\rangle^{\otimes m}) \end{aligned}$$

where we used the partial cyclicity property

$$\text{tr}_2 \left( (\mathbb{1} \otimes X_2) \gamma_{12} \right) = \text{tr}_2 \left( \gamma_{12} (\mathbb{1} \otimes X_2) \right)$$

in the last equality. We define

$$\sqrt{p(\varphi)} |e_\varphi\rangle = \binom{m+d-1}{m}^{1/2} (\mathbb{1}_k \otimes \langle\varphi|^{\otimes m}) |\psi\rangle_{A_1 \dots A_n} \in (\mathbb{C}^d)^{\otimes k}$$

where  $p(\varphi) \geq 0$  ensures that  $\langle e_\varphi | e_\varphi \rangle = 1$ .

Note that  $p(\varphi)$  is a probability density, i.e.,  $\int d\varphi p(\varphi) = 1$ .

Hence, we have  $\chi_{A_1 \dots A_k} = \int d\varphi p(\varphi) |e_\varphi\rangle\langle e_\varphi|$ .

To show:  $\int d\varphi p(\varphi) |e_\varphi\rangle\langle e_\varphi| \approx \int d\varphi p(\varphi) |\varphi\rangle\langle\varphi|^{\otimes k}$ .

First, we compute the average (squared) fidelity of  $|e_\varphi\rangle$  and  $|\varphi\rangle^{\otimes k}$ :

$$\int d\varphi \rho(\varphi) F(|e_\varphi\rangle, |\varphi\rangle^{\otimes k})^2 \propto (\mathbb{1}_k \otimes \langle \varphi |^{\otimes m}) | \varphi \rangle$$

$$= \int d\varphi \rho(\varphi) \langle e_\varphi | \varphi^{\otimes k} | e_\varphi \rangle$$

$$= \binom{m+d-1}{m} \int d\varphi \langle \varphi | \varphi^{\otimes k+m} | \varphi \rangle$$

$$= \binom{m+d-1}{m} \cdot \binom{m+d-1}{m}^{-1} \overbrace{\langle \varphi | \mathbb{1}_{k+m} | \varphi \rangle}^{=1}$$

$$= \binom{m+d-1}{m} \cdot \binom{k+m+d-1}{k+m}^{-1}$$

$$= \frac{(m+d-1)!}{m! (d-1)!} \frac{(k+m)! (d-1)!}{(k+m+d-1)!}$$

$$= \frac{(m+d-1) \dots (m+1)}{(k+m+d-1) \dots (k+m+1)}$$

$$\geq \left( \frac{m+1}{k+m+1} \right)^{d-1} = \left( 1 - \frac{k}{k+m+1} \right)^{d-1}$$

$$\geq 1 - \frac{k(d-1)}{k+m+1} \geq 1 - \frac{kd}{m}$$

Now we can finish the proof. Recall the

Fuchs-van-de-Graaf inequality  $D(\rho, \sigma) \leq \sqrt{1 - F(\rho, \sigma)^2}$

(which is actually an equality for pure states, see Ex.).

$$D(\varphi_{A_1 \dots A_n}, \int d\varphi \rho(\varphi) |\varphi \chi \varphi\rangle^{\otimes n})$$

$$= D\left(\int d\varphi \rho(\varphi) |e_\varphi \chi e_\varphi\rangle, \int d\varphi \rho(\varphi) |\varphi \chi \varphi\rangle^{\otimes n}\right)$$

convexity  
of norms  $\leq \int d\varphi \rho(\varphi) D(|e_\varphi \chi e_\varphi\rangle, |\varphi \chi \varphi\rangle^{\otimes n})$

$$\leq \int d\varphi \rho(\varphi) \sqrt{1 - F(e_\varphi, \varphi^{\otimes n})^2}$$

Jensen's  
ineq.  $\leq \left(\int d\varphi \rho(\varphi) (1 - F(e_\varphi, \varphi^{\otimes n})^2)\right)^{1/2}$

$$= \left(1 - \int d\varphi \rho(\varphi) F(e_\varphi, \varphi^{\otimes n})^2\right)^{1/2}$$

$$\leq \sqrt{1 - \left(1 - \frac{kd}{m}\right)} = \sqrt{\frac{kd}{m}}. \quad \square$$

### § 7.3 Extension to permutation-invariant mixed states

A state  $\rho_{A_1 \dots A_n}$  is called **permutation-invariant** if

$$Q_\pi \rho_{A_1 \dots A_n} Q_\pi^\dagger = \rho_{A_1 \dots A_n} \quad \text{for all } \pi \in S_n.$$

$$(Q_\pi \equiv \varphi(\pi))$$

Goal: prove de Finetti theorem for permutation-invariant  $\rho$ .

Strategy: use our theorem from §7.2 for pure states!

Need to relate perm.-inv. states to pure states in  $\text{Sym}^n(\mathbb{C}^d)$

**Lem** Let  $\mathcal{H}_{A_i} = \mathbb{C}^d$  for  $i=1, \dots, n$  and  $\rho_{A_1 \dots A_n}$  be permutation-invariant. Then  $\rho_{A_1 \dots A_n}$  has a purification

$$|\psi^S\rangle \in \text{Sym}^n(\mathbb{C}^d \otimes \mathbb{C}^d).$$

Proof: Let  $\rho_{A_1 \dots A_n} = \sum_{\lambda \in \text{spec}(\rho)} \lambda P_\lambda$  be a spectral

decomposition, where  $\text{spec}(\rho)$  is the set of distinct eigenvalues

of  $\rho$  with corresponding orthogonal projector  $P_\lambda$  onto the

eigenspace  $\mathcal{H}_\lambda$ . Since  $\rho = Q_\pi \rho Q_\pi^\dagger$  for all  $\pi \in S_n$ ,

we have for any  $\lambda \in \text{spec}(\rho)$  and  $|\phi\rangle \in \mathcal{H}_\lambda$  that

$$\lambda |\phi\rangle = \rho |\phi\rangle = Q_\pi \rho Q_\pi^\dagger |\phi\rangle$$

and hence  $Q_\pi^\dagger |\phi\rangle \in \mathcal{H}_\lambda$  for all  $\pi \in S_n$ , i.e., the eigenspaces

are permutation-invariant too, and  $P_\lambda Q_\pi = Q_\pi P_\lambda$  for all

$\pi \in S_n$ ,  $\lambda \in \text{spec}(\rho)$ . Define  $M = \sum_{\lambda \in \text{spec}(\rho)} \sqrt{\lambda} P_\lambda$ ,

then clearly also  $Q_\pi M = M Q_\pi$  for all  $\pi \in S_n$ .

Let now  $|\phi\rangle_{A_1 \dots A_n R_1 \dots R_n} := \sum_{x=1}^{d^n} |x\rangle_{A^n} \otimes |x\rangle_{R^n}$ , where  $\{|x\rangle\}_{x=1}^{d^n}$  is a product basis for  $(\mathbb{C}^d)^{\otimes n}$ , and set

$$|\psi^S\rangle = (M \otimes \mathbb{1}_{R^n}) |\phi\rangle.$$

Then:  $\text{tr}_{R^n} \mathcal{Z}_{A^n R^n}^S = M \underbrace{(\text{tr}_{R^n} \phi_{A^n R^n})}_{= \mathbb{1}_{A^n}} M^\dagger$

$$= M M^\dagger = \sum_{\lambda, \lambda'} \sqrt{\lambda \lambda'} \underbrace{P_\lambda P_{\lambda'}}_{= \delta_{\lambda \lambda'} P_\lambda} = S_{A^n}$$

$\rightarrow$  Note: If  $S_n$  acts on both  $A_1 \dots A_n$  and  $R_1 \dots R_n$  via  $Q_\pi$ , then  $S_n$  acts on  $A_1 \dots A_n R_1 \dots R_n \cong A_1 R_1 \dots A_n R_n$  via  $Q_\pi \otimes Q_\pi$ .

$$\rightarrow (Q_\pi \otimes Q_\pi) |\psi^S\rangle = (Q_\pi \otimes Q_\pi) (M \otimes \mathbb{1}) |\phi\rangle$$

$$= (Q_\pi M \otimes \mathbb{1}) (\mathbb{1} \otimes Q_\pi) |\phi\rangle$$

transpose trick  $\rightarrow$

$$= (Q_\pi M Q_\pi^\dagger \otimes \mathbb{1}) |\phi\rangle$$

$[M, Q_\pi] = 0$   $\rightarrow$

$$= (M Q_\pi Q_\pi^\dagger \otimes \mathbb{1}) |\phi\rangle$$

$Q_\pi^\dagger = Q_\pi^\top$   $\rightarrow$

$$= (M \otimes \mathbb{1}) |\phi\rangle$$

$$= |\psi^S\rangle \quad \text{for all } \pi \in S_n$$

$$\Rightarrow |\psi^S\rangle \in \text{Sym}^n(\mathbb{C}^d \otimes \mathbb{C}^d). \quad \square$$



**Prop** Let  $\mathcal{H}_{A_i} = \mathbb{C}^d$  for  $i=1, \dots, n$  and  $\rho_{A_1 \dots A_n}$  be a permutation-invariant state. Then for any  $k < n$ ,

$$D(\rho_{A_1 \dots A_k}, \int d\mu(\sigma) \sigma^{\otimes k}) \leq \sqrt{\frac{d^2 k}{n-k}},$$

where  $d\mu(\sigma)$  is a measure on the space of mixed states on  $\mathbb{C}^d$  that depends on  $\rho$ .

Proof: Let  $|\psi^S\rangle_{A_1 R_1 \dots A_k R_k} \in \text{Sym}^n(\mathbb{C}^d \otimes \mathbb{C}^d)$  be a symmetric purification of  $\rho$ . Applying the sym. subspace de Finetti theorem shows that

$$D\left(\psi^S_{A_1 R_1 \dots A_k R_k}, \int d\psi p(\psi) |\psi\rangle\langle\psi|_{AR}^{\otimes k}\right) \leq \sqrt{\frac{d^2 k}{n-k}}$$

for a suitable prob. density  $p(\psi)$ . The claim now follows from the monotonicity of  $D(\cdot, \cdot)$  under partial trace:

$$\begin{aligned} & D(\rho_{A_1 \dots A_k}, \int d\psi p(\psi) \text{tr}_R |\psi\rangle\langle\psi|_{AR}^{\otimes k}) \\ & \leq D\left(\psi^S_{A_1 R_1 \dots A_k R_k}, \int d\psi p(\psi) |\psi\rangle\langle\psi|_{AR}^{\otimes k}\right) \\ & \leq \sqrt{\frac{d^2 k}{n-k}}. \end{aligned}$$

□